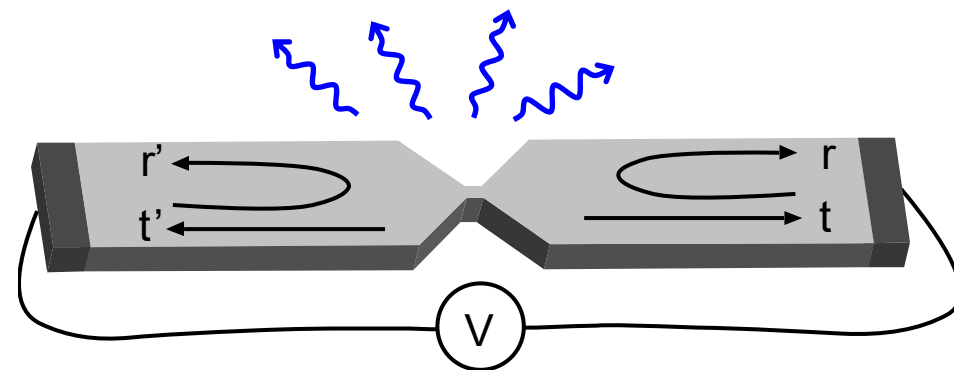


Josephson Photonics - Quantum Optics and Cooper Pair Tunneling

Joachim Ankerhold

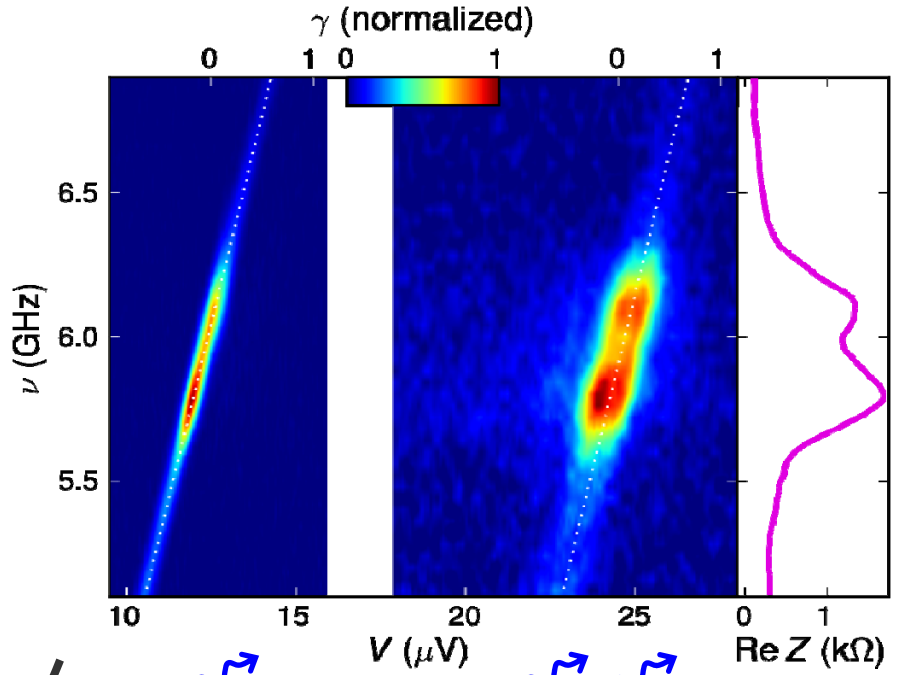
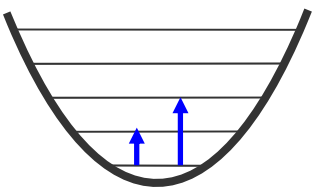
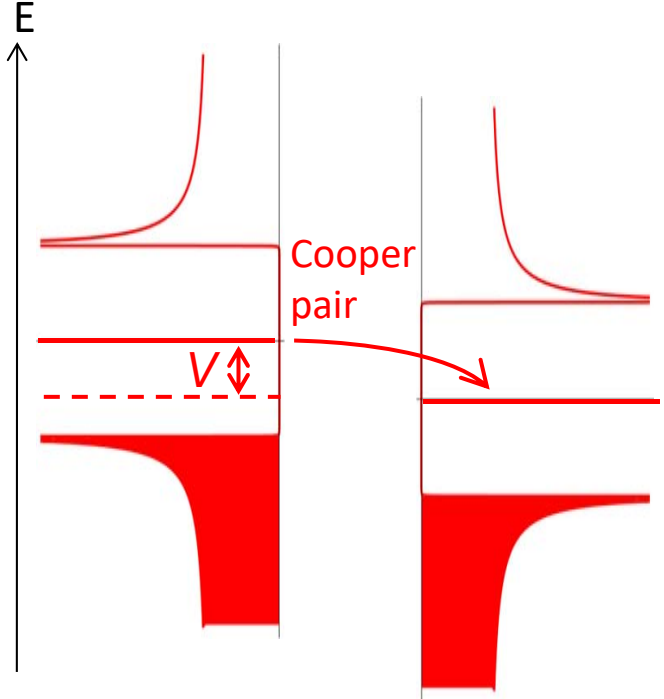
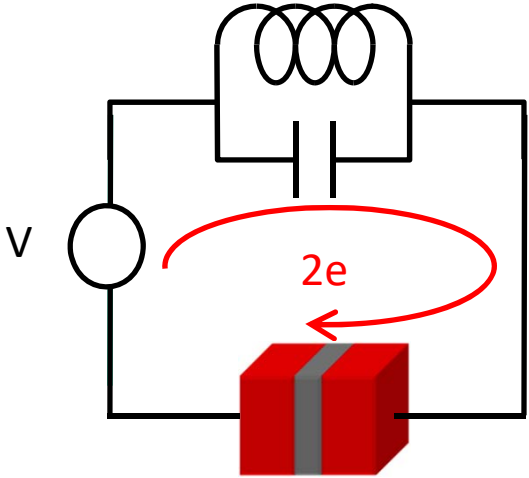
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Photon emission from a dc-voltage biased Josephson junction



$2eV = \hbar\omega$

$2eV = 2\hbar\omega$

Dynamical Coulomb blockade

$$H = \hbar\omega a^\dagger a - E_J \cos[2eVt/\hbar + \varphi]$$

$$\varphi = \sqrt{r} (a^\dagger + a)$$

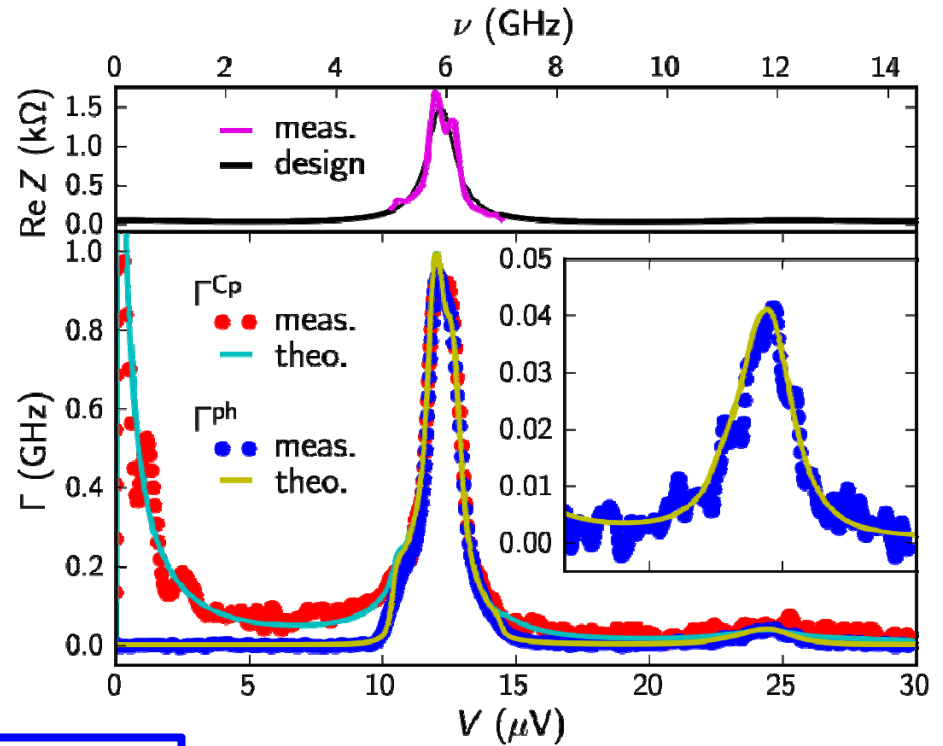
Charge-light coupling:
„fine structure constant“

$$r = \pi \frac{Z_{\text{res}}}{R_Q}$$

$$\begin{aligned} \Gamma_{\rightarrow}^{2e}(V) &= \frac{\pi E_J^2}{2\hbar} \sum_p |\langle p | e^{i\varphi} | 0 \rangle|^2 \delta(2eV - p\hbar\omega) \\ &= \frac{\pi E_J^2}{2h} \sum_p \frac{\exp(-r) r^p}{p!} \delta(2eV - p\hbar\omega) \end{aligned}$$

$$\Gamma_{\rightarrow}^{h\nu}(V = p\hbar\omega/2e) = p \Gamma_{\rightarrow}^{2e}$$

$$\Gamma_{\rightarrow}^{2e} \ll \gamma$$



From Coulomb blockade to nonlinear quantum oscillator

Rotating frame $\omega_J \equiv 2eV/\hbar \approx p\omega$

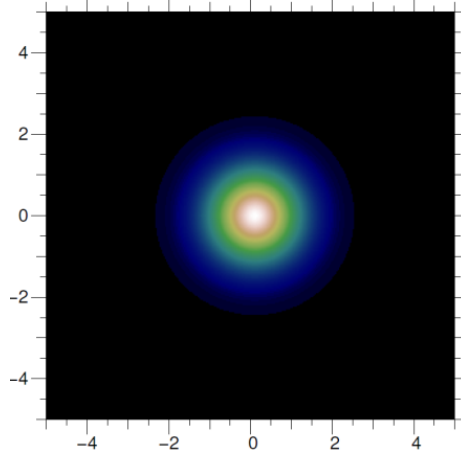
$$H_p = \hbar\Delta_p n + (-i)^p \frac{E_J^*}{2} : [(a^\dagger)^p + (-1)^p a^p] \frac{J_p(2\sqrt{r n})}{n^{p/2}} : , \quad E_J^* = E_J e^{-r/2}$$

$$\dot{\rho}(t) = -\frac{i}{\hbar} [H_p, \rho(t)] + \frac{\gamma}{2} [2a\rho(t)a^\dagger - a^\dagger a\rho(t) - \rho(t)a^\dagger a] \longrightarrow \dot{\rho}(t) = 0$$

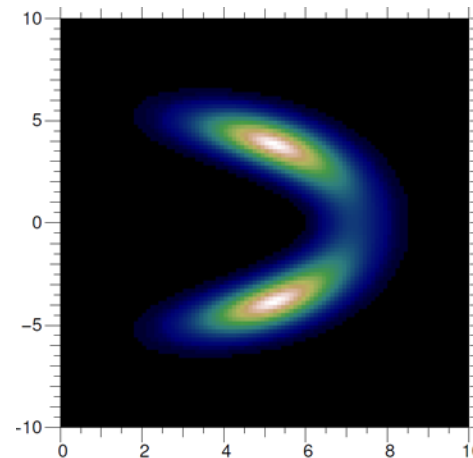
Cavity field:

$$\Gamma_{\rightarrow}^{2e} \ll \gamma$$

$$p = 1$$



$$\Gamma_{\rightarrow}^{2e} \gg \gamma$$



Josephson photonics

PHYSICAL REVIEW B **90**, 020506(R) (2014)



Realization of a single-Cooper-pair Josephson laser

Fei Chen,¹ Juliang Li,¹ A. D. Armour,² E. Brahim, ¹ Joel Stettenheim,¹ A. J. Sirois,³ R. W. Simmonds,⁴
M. P. Blencowe,¹ and A. J. Rimberg^{1,*}

PRL **119**, 137001 (2017)

PHYSICAL REVIEW LETTERS

week ending
29 SEPTEMBER 2017

REPORT

Science 355, 939 (2017)

OPTOELECTRONICS

Demonstration of an ac Josephson junction laser

M. C. Cassidy,¹ A. Bruno,¹ S. Rubbert,² M. Irfan,² J. Kammhuber,¹ R. N. Schouten,^{1,2}
A. R. Akhmerov,² L. P. Kouwenhoven^{1,2,*}

Leppäkangas et al, PRL 110, 267004 (2013)
Armour et al, PRL 111, 247001 (2013)
Altimiras et al, PRL 112, 236803 (2014)
S. Meister et al, PRB 92, 174532 (2015)

ARTICLE

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DOI: 10.1038/ncomms6562

OPEN

Photon-assisted tunnelling with nonclassical light

J.-R. Souquet^{1,2}, M.J. Woolley³, J. Gabelli¹, P. Simon¹ & A.A. Clerk²

Emission of Nonclassical Radiation by Inelastic Cooper Pair Tunneling

M. Westig,¹ B. Kubala,² O. Parlavecchio,¹ Y. Mukharsky,¹ C. Altimiras,¹ P. Joyez,¹ D. Vion,¹ P. Roche,¹ D. Esteve,¹
M. Hofheinz,^{1,*} M. Trif,³ P. Simon,³ J. Ankerhold,^{2,†} and F. Portier^{1,‡}

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PHYSICAL REVIEW B **96**, 214509 (2017)

Noise switching at a dynamical critical point in a cavity-conductor hybrid

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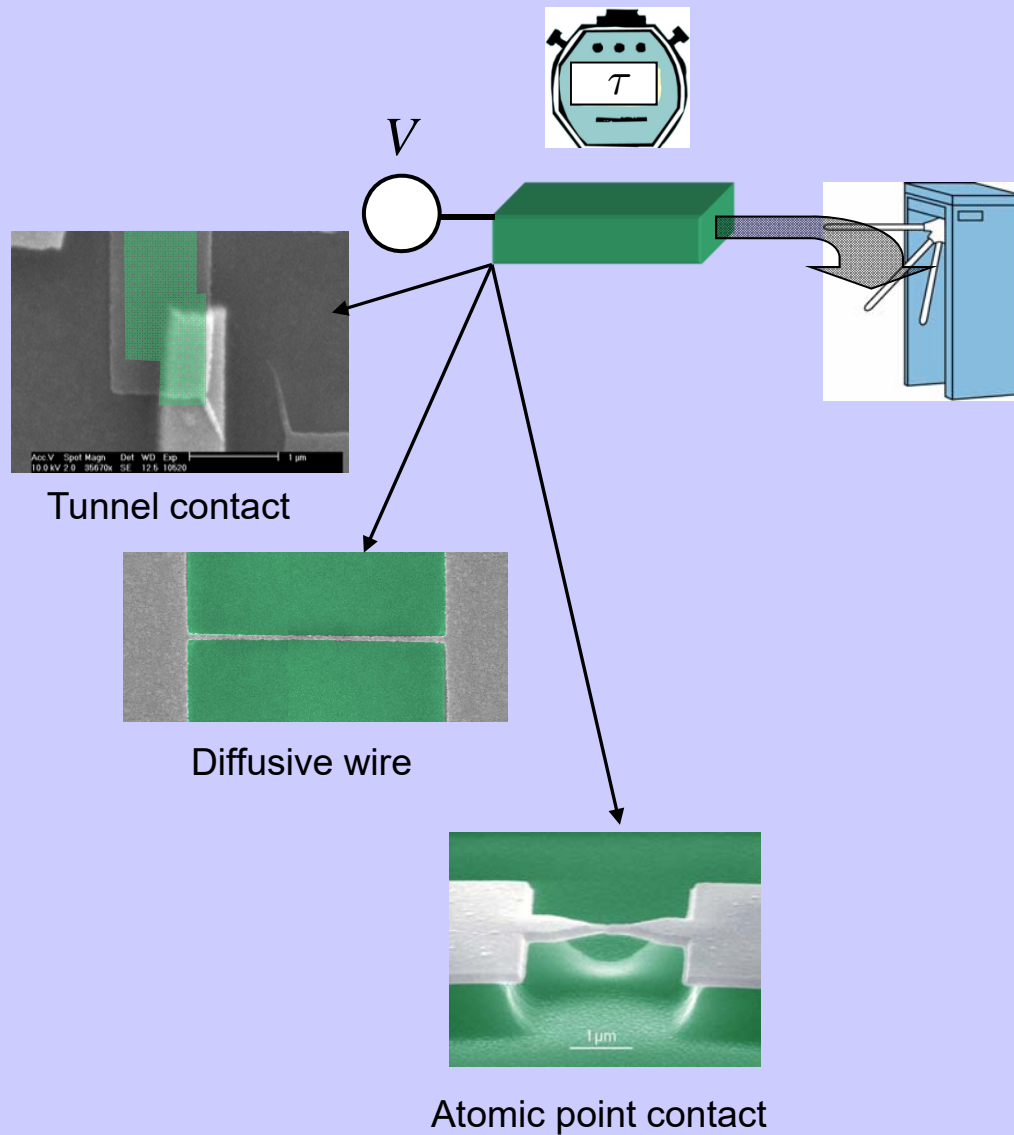
Gramich et al, PRL 113, 027001 (2014)

Armour et al, PRB 91, 184508 (2015)

Ast et al, Nature Comm. 7, 13009 (2016)

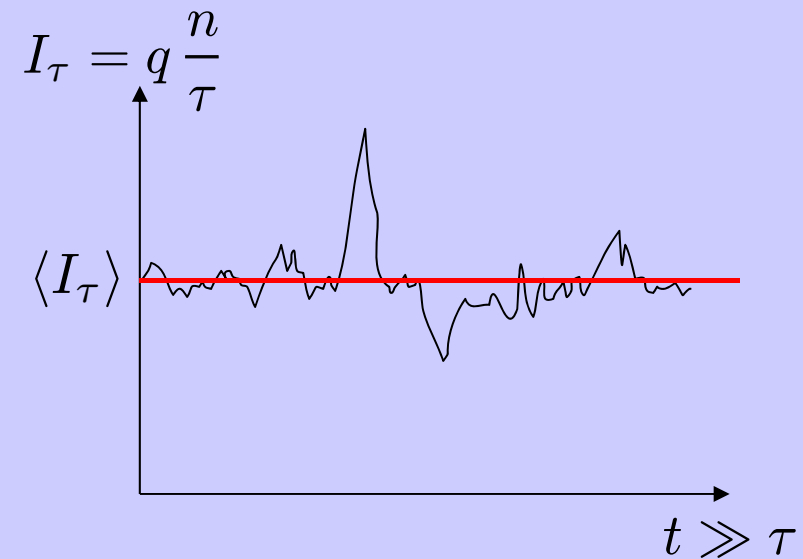
Dykman et al, arXiv: 1702.07931 (2017)

Counting charges

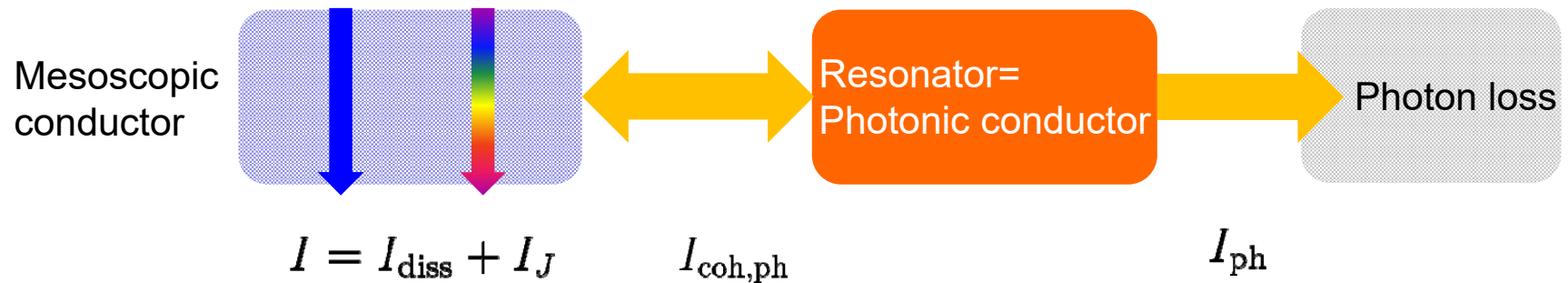


R. Landauer:

The noise is the signal



Counting charges via photonic conductor



Steady state $\langle I_J \rangle / 2e = \langle I_{\text{coh,ph}} \rangle = \langle I_{\text{ph}} \rangle$

Josephson @ low voltage $I_{\text{diss}} = 0$

Cooper pair noise from photon noise ?

Photon counting statistics

Generating function $\exp[\mathcal{F}_{ph}(\chi, t)] = \text{Tr}\{\exp[\mathcal{L}_\chi t] \rho(0)\}$

$$\mathcal{L}_\chi = \mathcal{L}_p + \frac{\gamma}{2}(2e^{-i\chi} a \rho a^\dagger - a^\dagger a \rho - \rho a^\dagger a)$$

Unitary transformation in Liouville space $\mathcal{L}_\chi \rightarrow \mathcal{L}_{p\chi}^{\text{coh,ph}}$

$$\mathcal{L}_{p\chi}^{\text{coh,ph}} \rho = \frac{i}{\hbar}(H_\chi \rho - \rho H_{-\chi}) + \mathcal{L}_\gamma$$
$$H_\chi = H_p(a \rightarrow e^{i\chi/2} a)$$

New generating function $\exp[\mathcal{F}_{\text{coh,ph}}(\chi, t)] = \text{Tr}\{\exp[\mathcal{L}_{p\chi}^{\text{coh,ph}} t] \rho(0)\}$

Coherent photon transfer – quasi-probability $\tilde{P}(N, t)$

Photon counting statistics

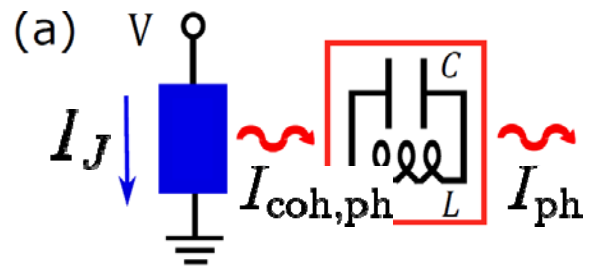
$$\mathcal{F}_{\text{coh,ph}}(p\chi, \gamma t \gg 1) = \mathcal{F}_{\text{ph}}(\chi, \gamma t \gg 1)$$

Cumulants photon numbers $\lim_{\gamma t \rightarrow \infty} \langle \langle N^k \rangle \rangle_{\text{coh,ph}} = \langle \langle N^k \rangle \rangle_{\text{ph}} / p^k$

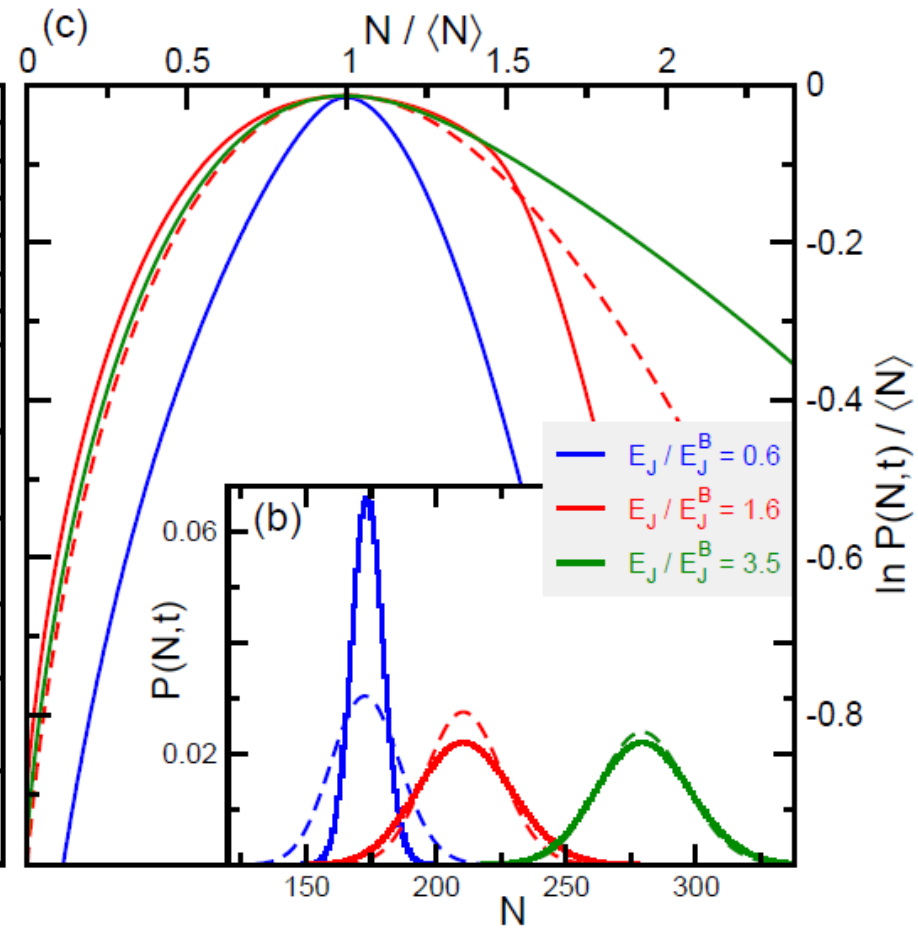
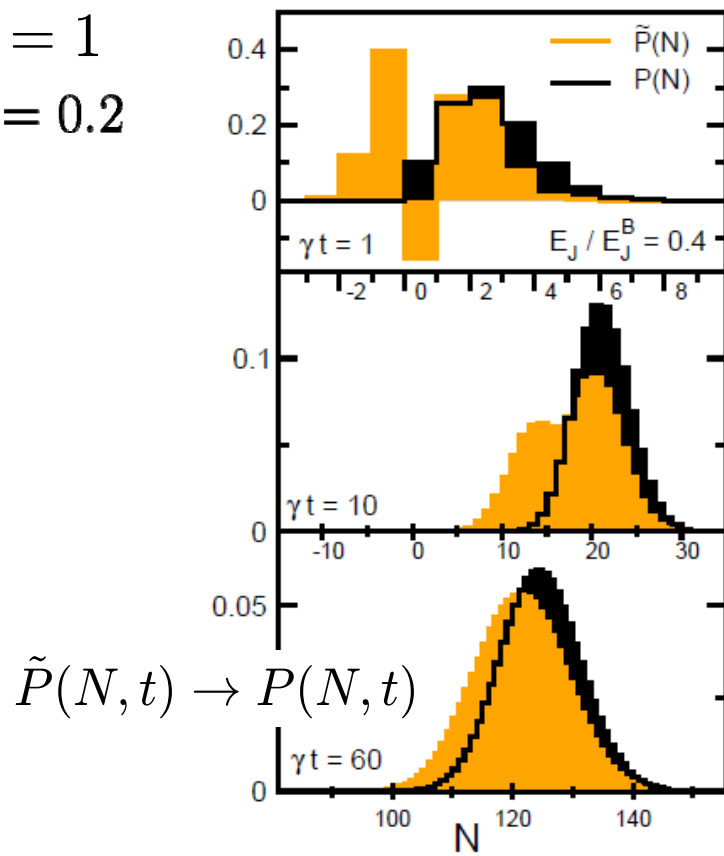
Long times:

Statistics of emitted photons = statistics of transferred Cooper pairs

Full counting statistics

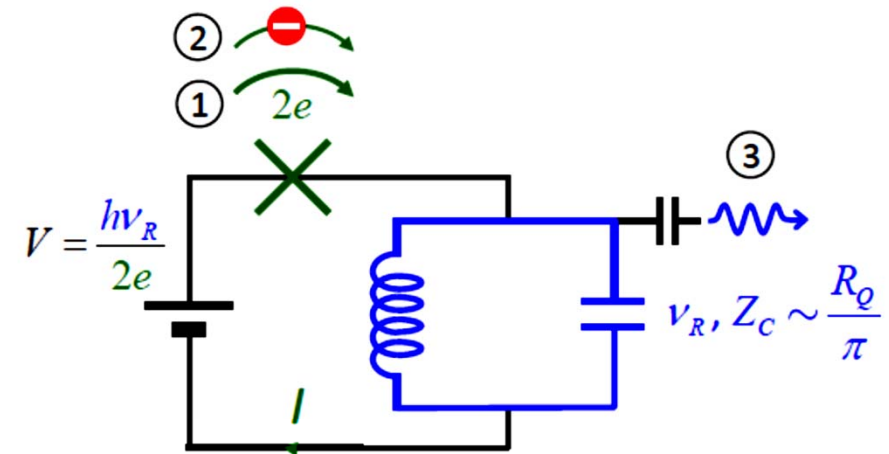
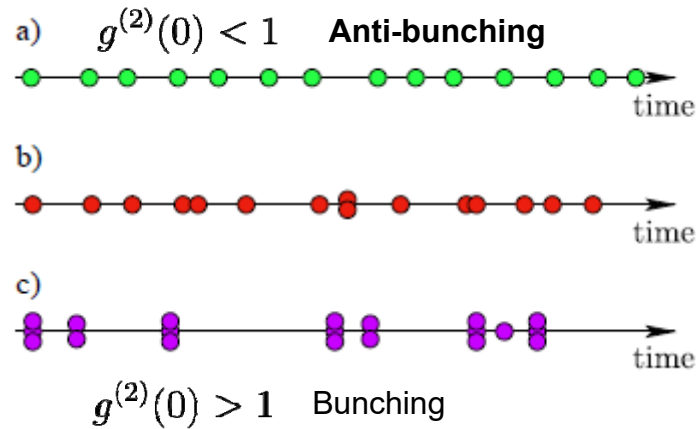


$p = 1$
 $r = 0.2$



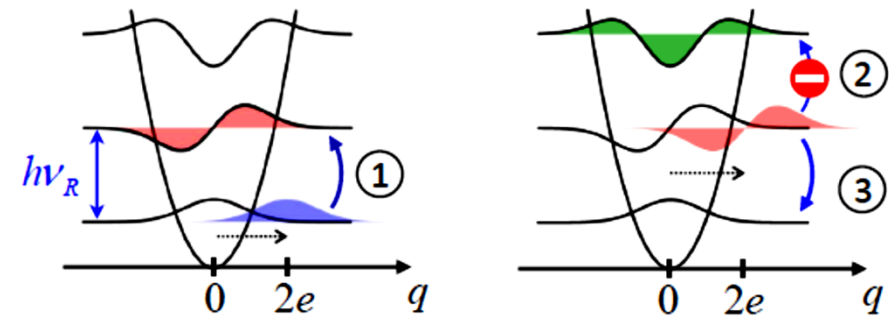
Designing photon statistics: Towards single photon emission

$$g^{(2)}(\tau) = \frac{\langle a^\dagger(\tau) a^\dagger a a(\tau) \rangle_{st}}{\langle n \rangle_{st}^2}$$

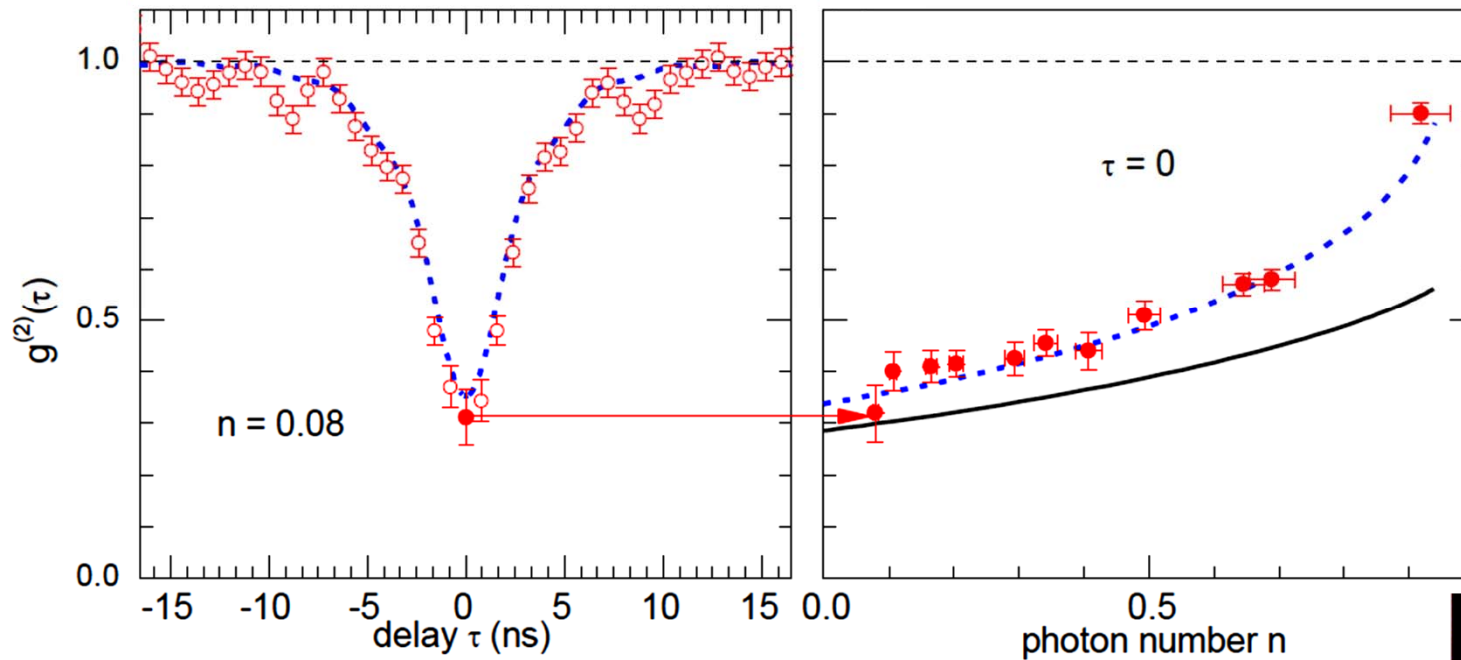


Weak driving

$$g^{(2)}(0) = (1 - r/2)^2$$



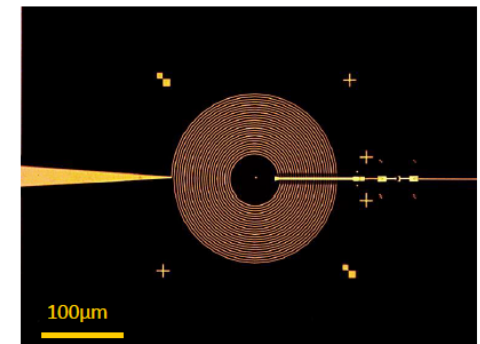
Towards single photon sources



$\sim 10^7$ photons/s

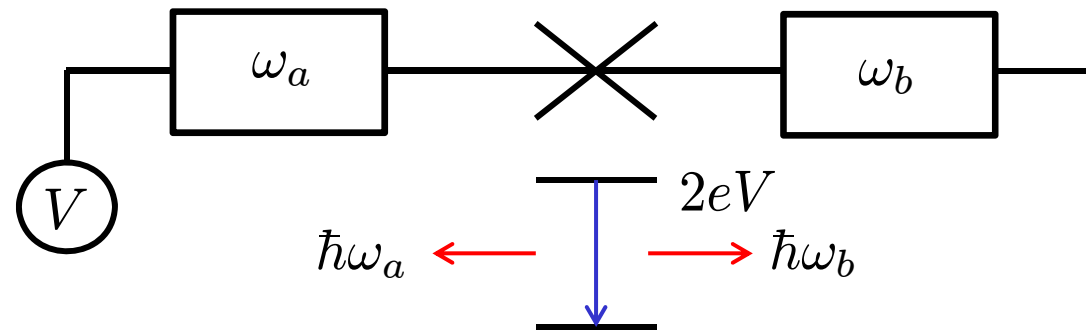
$$r \sim 1$$

$$Z_{\text{res}} = \sqrt{L/C}$$



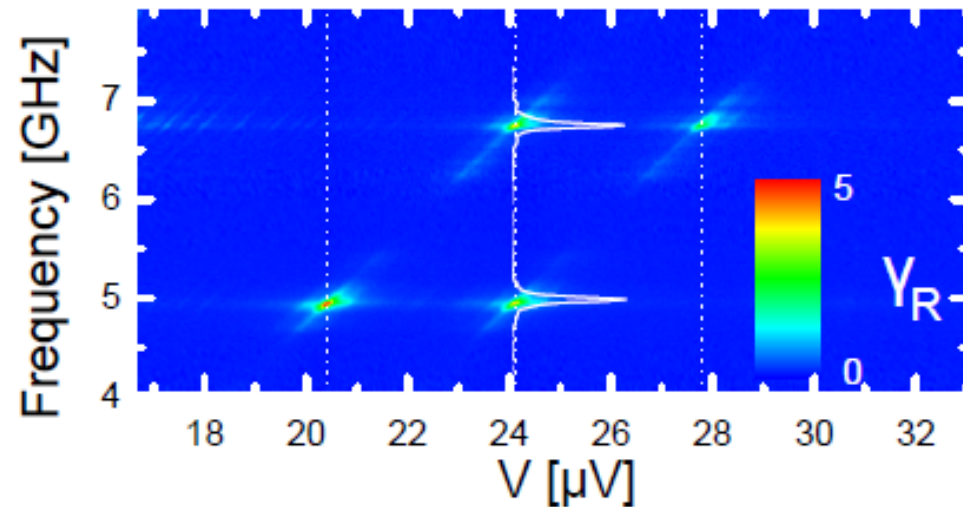
Designing photon statistics: Towards entangled photons

$$2eV = \hbar(\omega_a + \omega_b)$$

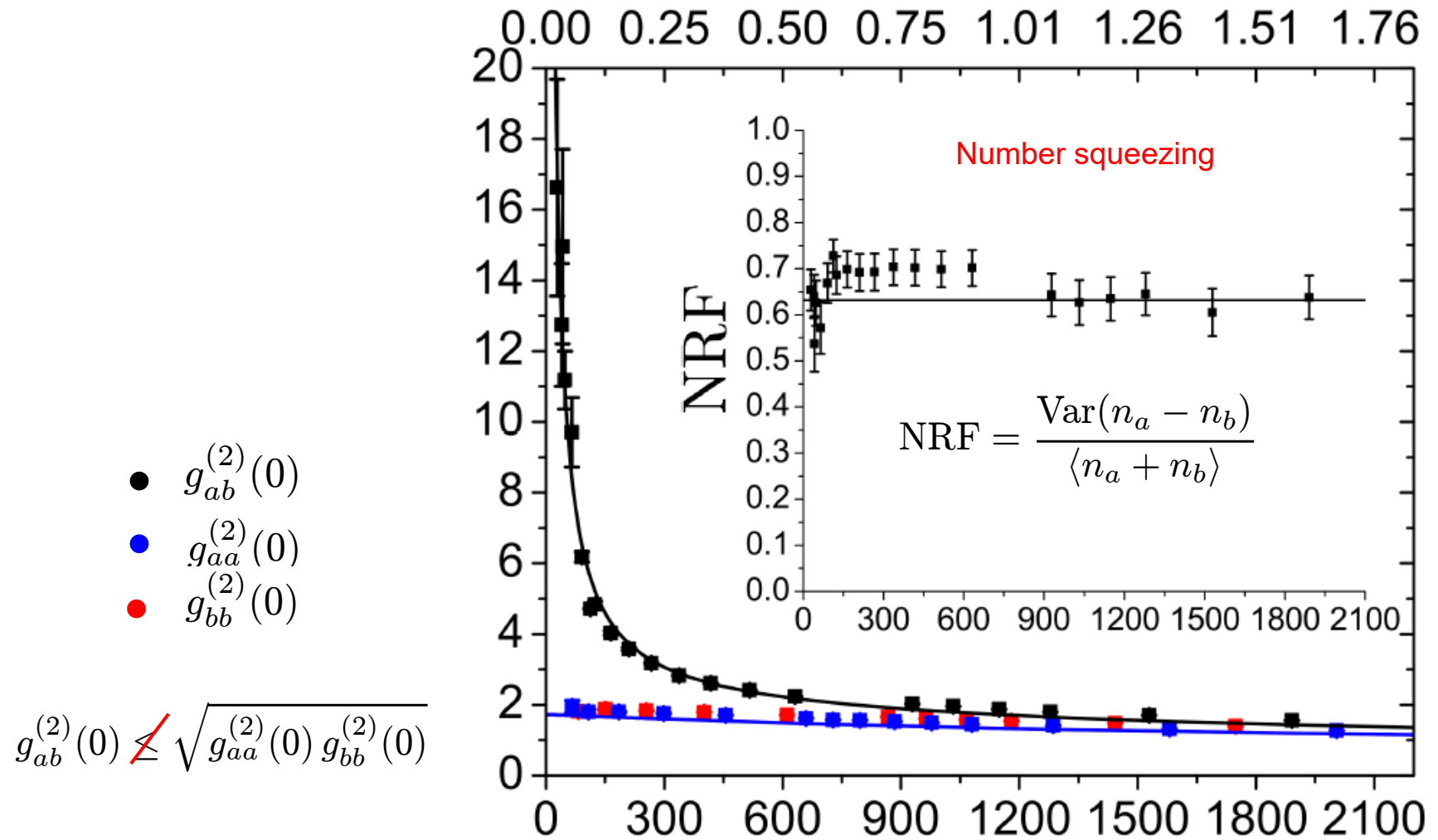


Non-classical light:
violation of Cauchy-Schwartz

$$g_{ab}^{(2)}(0) \not\leq \sqrt{g_{aa}^{(2)}(0) g_{bb}^{(2)}(0)}$$

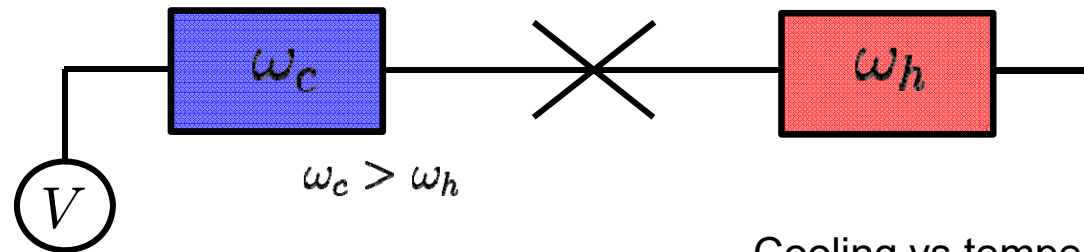


Two cavities: correlated photons

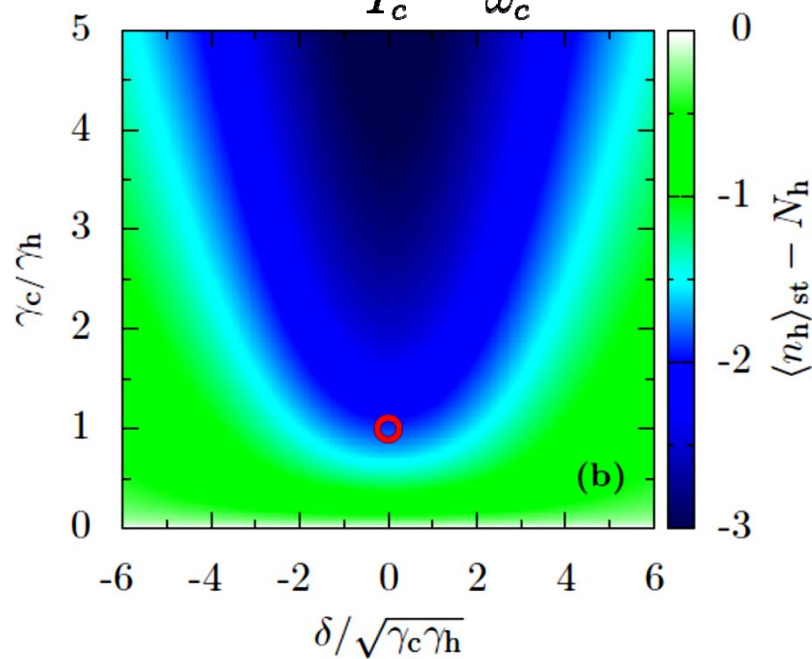


Thermodynamic device: refrigerator

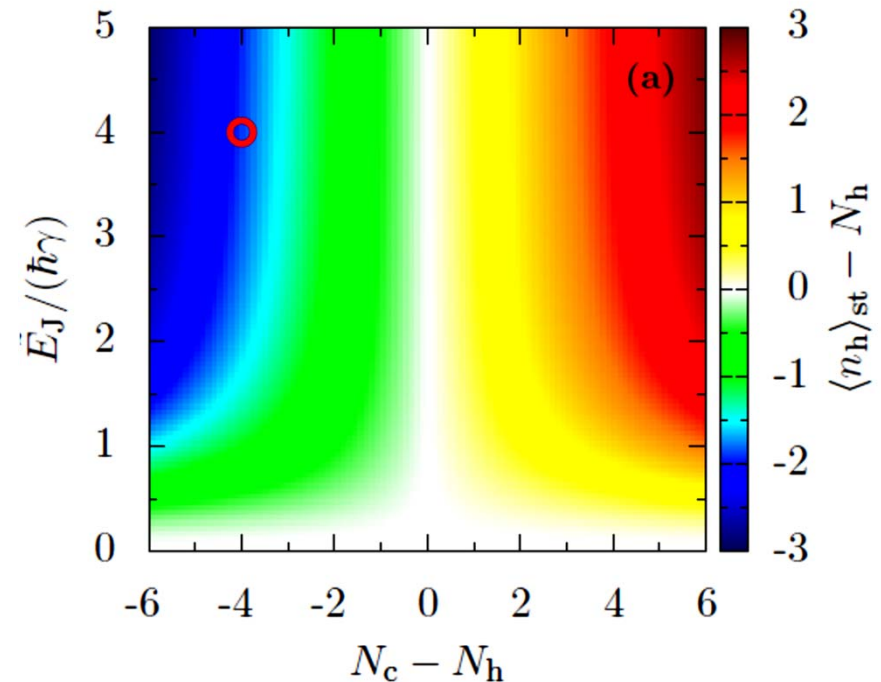
$$2eV + \hbar\omega_h = \hbar\omega_c$$



Cooling vs de-tuning $\frac{T_h}{T_c} < \frac{\omega_h}{\omega_c}$, ($\hbar\omega_c < \Delta$)

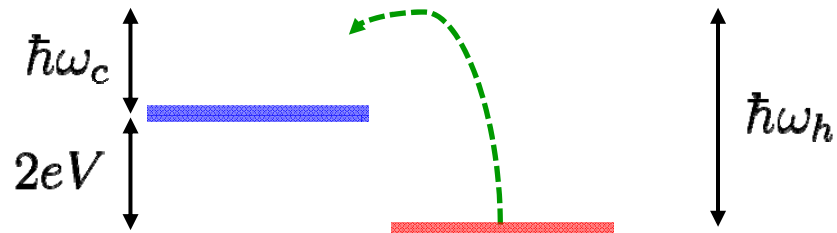
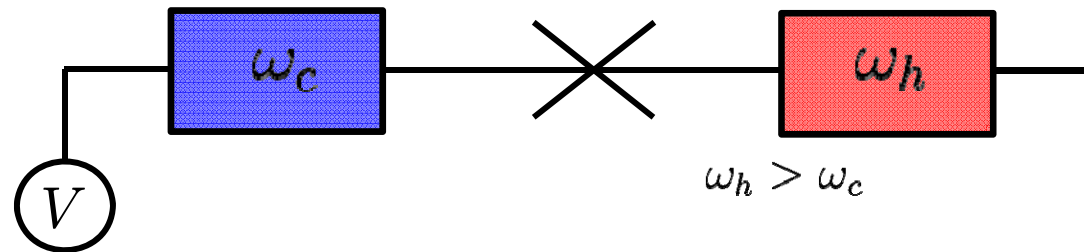


Cooling vs temperature gradient

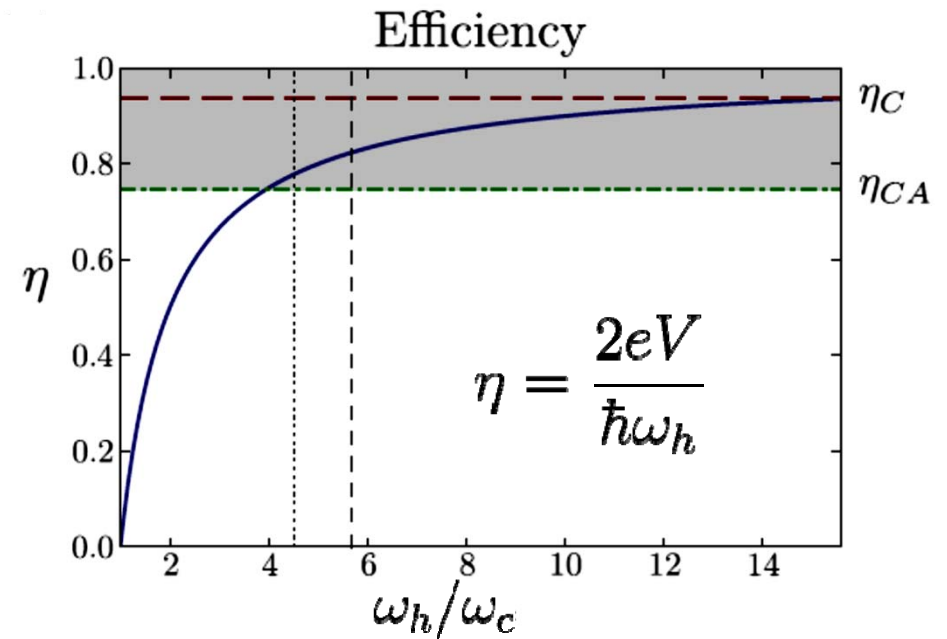


Thermodynamic device: Heat engine

$$2eV + \hbar\omega_c = \hbar\omega_h$$



Net current against bias $\frac{\omega_c}{\omega_h} > \frac{T_c}{T_h}$



- Quantum optics of mesoscopic conductors: Josephson photonics
- Cooper pair noise \longleftrightarrow emitted photon statistics
- Single photon emission
- Correlated photon pairs
- (Quantum) Thermodynamics devices

